

which is equivalent to Riemann's relation $g = \frac{1}{2} b_1(M)$.

$$\begin{aligned} \Gamma \quad C_1(M) &= 2(1-g) = \underbrace{1}_{\dim H^0(M)} - b_1(M) + \underbrace{1}_{\dim H^2(M)} = 2 - b_1(M) \\ \Rightarrow \quad g &= \frac{1}{2} b_1(M) \end{aligned} \quad \sqcup$$

This we proved by harmonic theory. Γ See p 124. \sqcup

For a surface, we have Noether's formula

$$\chi(\mathcal{O}_M) = \frac{C_1(M)^2 + C_2(M)}{12},$$

which we will prove and use extensively in the next chapter.

$$\Gamma \text{ Since } n=2, \quad Td_2 = \frac{P_1(A)^2 + P_2(A)}{12},$$

by Hirzebruch - Riemann - Roch formula,

$$Td_2 = \frac{C_1(M)^2 + C_2(M)}{12} = \chi(\mathcal{O}_M). \quad \sqcup$$

Unfortunately, our analogy between the Gauss - Bonnet III and Riemann - Roch formulas fails in one crucial aspect: while any differentiable manifold has many C^∞ vector fields to use as pros in the proof of Gauss - Bonnet III, relatively few compact manifolds have any global holomorphic vector fields. (Cf. the theorem of Carrell and Lieberman proved in Section 4 of Chapter 5.) Of course, since the Riemann - Roch formula itself has nothing to do with the vector field v used to obtain it,