

□

$$\begin{aligned} \|G(\psi)\|_{s+2}^2 &= \sum_{z \neq 0} (1 + \|z\|^2)^{s+2} \left| \frac{\psi_z}{\|z\|^2} \right|^2 \\ &= \sum_{z \neq 0} \frac{(1 + \|z\|^2)^{s+2}}{\|z\|^4} |\psi_z|^2 = \sum_{z \neq 0} \frac{(1 + \|z\|^2)^2}{\|z\|^4} (1 + \|z\|^2)^s |\psi_z|^2 \end{aligned}$$

$$< \sum_{z \neq 0} 2 (1 + \|z\|^2)^s |\psi_z|^2 \leq 2 \sum_{z \in \mathbb{Z}^n} (1 + \|z\|^2)^s |\psi_z|^2 = 2 \|\psi\|_s^2 \quad \text{---}$$

$\Rightarrow G$ is bounded

In case, ψ is perpendicular to the harmonic space, ^(weak solution space of $\Delta \varphi = 0$)

$\varphi = G(\psi)$ gives a weak solution to (*). By the Sobolev lemma, if $\psi \in C^\infty(T)$, then $\varphi \in C^\infty(T)$ and φ is a solution of (*) in the usual sense.

□ $G(\psi) = \varphi \in H_{s+2}(T)$ $\varphi \in \bigcap H_s = H_\infty$ □

Finally, by the Rellich lemma,

$G : L^2(T) \longrightarrow L^2(T)$ is a compact, self-adjoint operator.

$$\square \quad G : L^2(T) = H_0 \xrightarrow{\text{bounded}} H_2(T) \xleftarrow{\text{compact}} H_0(T) = L^2(T)$$

$$\langle G(\psi), \phi \rangle \stackrel{?}{=} \langle \psi, G(\phi) \rangle, \psi, \phi \in L^2(T)$$

$$\psi = \sum \psi_z e^{\bar{z}\langle z, x \rangle}, \quad \phi = \sum \phi_z e^{\bar{z}\langle z, x \rangle}$$

$$G(\psi) = \sum_{z \neq 0} \psi_z \frac{e^{\bar{z}\langle z, x \rangle}}{\|z\|^2} \quad G(\phi) = \sum_{z \neq 0} \frac{\phi_z}{\|z\|^2} e^{\bar{z}\langle z, x \rangle}$$

$$\Rightarrow \langle G(\psi), \phi \rangle = \left\langle \sum_{z \neq 0} \psi_z \frac{e^{\bar{z}\langle z, x \rangle}}{\|z\|^2}, \sum_{z' \neq 0} \phi_{z'} e^{\bar{z}'\langle z', x \rangle} \right\rangle$$

$$= \sum_{z \neq 0} \psi_z \bar{\phi}_z \frac{1}{\|z\|^2} \stackrel{?}{=} \langle \psi, G(\phi) \rangle$$

$$= \left\langle \sum_z \psi_z e^{\bar{z}\langle z, x \rangle}, \sum_{z' \neq 0} \frac{\phi_{z'}}{\|z'\|^2} e^{\bar{z}'\langle z', x \rangle} \right\rangle = \sum_{z \neq 0} \frac{\psi_z \bar{\phi}_z}{\|z\|^2} \quad \square$$