

Finally, the $\bar{\partial}$ -Laplacian on E is defined by

$$\Delta = \bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial} : A^{p,q}(E) \longrightarrow A^{p,q}(E).$$

An E -valued form φ is called harmonic if $\Delta \varphi = 0$.
(Again, harmonic forms φ are exactly the forms of smallest norm in their Dolbeault cohomology class $\varphi + \bar{\partial} A^{p,q-1}(E)$.) We let

$$\mathcal{H}^{p,q}(E) = \text{Ker } \Delta \quad \text{be the harmonic space.}$$

Now, the analytic part of the proof of the Hodge theorem for the $\bar{\partial}$ -Laplacian on ordinary differential forms on M is essentially local: we can always find appropriate solutions of $\Delta \varphi = 0$ in the completion of $A^{p,q}(M)$ in the L_2 -norm; the problem is to show that these solutions are in fact C^∞ . Writing out E -valued forms in terms of a frame for E , all the local estimates used in the proof of the Hodge theorem for $A^*(M)$ go over to $A^{p,q}(E)$ — the only difference is that in each estimate we will get lower-order terms involving the coefficient functions for the metric on E as well as the metric on T^*M , and these can be estimated out as before. Thus the Hodge theorem holds for the $\bar{\partial}$ -Laplacian on E , that is:

1. $\mathcal{H}^{p,q}(E)$ is finite dimensional, and
2. If \mathcal{H} denotes the orthogonal projection $A^{p,q}(E) \rightarrow \mathcal{H}^{p,q}(E)$, there exists an operator

$$G : A^{p,q}(E) \longrightarrow A^{p,q}(E)$$

$$\text{s.t.} \quad G(\mathcal{H}^{p,q}(E)) = 0,$$