

But at  $t=0$

$$\frac{\partial}{\partial t} (\psi + t \bar{\partial} \eta, \psi + t \bar{\partial} \eta) = 2 \operatorname{Re}(\psi, \bar{\partial} \eta)$$

$$\text{and } \frac{\partial}{\partial t} (\psi + t \bar{\partial}(\bar{\partial} \eta), \psi + t \bar{\partial}(\bar{\partial} \eta)) = 2 \operatorname{Im}(\psi, \bar{\partial} \eta)$$

$$\Rightarrow \operatorname{Re}(\psi, \bar{\partial} \eta) = 0 \text{ \& } \operatorname{Im}(\psi, \bar{\partial} \eta) = 0 \Rightarrow (\psi, \bar{\partial} \eta) = 0$$

"  $(\bar{\partial}^* \psi, \eta)$

$$\text{for all } \eta \in A^{p,q-1}(M) \Rightarrow \bar{\partial}^* \psi = 0. \quad \text{Q.E.D.}$$

So, at least formally, the ~~Deo~~ Dolbeault cohomology group  $H_{\bar{\partial}}^{p,q}(M) = Z_{\bar{\partial}}^{p,q}(M) / \bar{\partial} A^{p,q-1}(M)$  is represented exactly by the solutions of the two first-order equations

$$\bar{\partial} \psi = 0, \quad \bar{\partial}^* \psi = 0.$$

These two may be replaced by the single-order equation

$$\Delta_{\bar{\partial}} \psi = (\bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}) \psi = 0;$$

clearly  $\bar{\partial} \psi = \bar{\partial}^* \psi = 0 \Rightarrow \Delta \psi = 0$ , and conversely,

$$(\Delta_{\bar{\partial}} \psi, \psi) = \langle \bar{\partial} \bar{\partial}^* \psi, \psi \rangle + \langle \bar{\partial}^* \bar{\partial} \psi, \psi \rangle = \|\bar{\partial}^* \psi\|^2 + \|\bar{\partial} \psi\|^2$$

$$\text{so } \Delta_{\bar{\partial}} \psi = 0 \Rightarrow \bar{\partial} \psi = 0 = \bar{\partial}^* \psi.$$

The operator

$$\Delta_{\bar{\partial}} : A^{p,q}(M) \longrightarrow A^{p,q}(M) \quad \text{is called the } \bar{\partial}\text{-Laplacian,}$$

or simply the Laplacian (written  $\Delta$ ) if no ambiguity is likely.

Differential forms satisfying the Laplacian equation  $\Delta \psi = 0$  are called harmonic forms; the space of harmonic forms of type  $(p,q)$  is denoted  $\mathcal{H}^{p,q}(M)$  and called the harmonic space. What the above formal argument suggests