

$$\nabla \psi = \nabla' \psi + \nabla'' \psi.$$

$$\|\nabla \psi\|^2 = \|\nabla' \psi\|^2 + \|\nabla'' \psi\|^2 \quad \text{since} \quad \langle \varphi_i, \bar{\varphi}_j \rangle = 0.$$

Remark 1. From $\langle \Delta \psi, \psi \rangle = \|\bar{\nabla} \psi\|^2 + \langle A' \psi, \psi \rangle$,

$$\langle \Delta \psi, \psi \rangle - \|\bar{\nabla} \psi\|^2 = \langle A' \psi, \psi \rangle$$

$$\Rightarrow 2|\langle \Delta \psi, \psi \rangle - \|\bar{\nabla} \psi\|^2| = 2|\langle A' \psi, \psi \rangle| \leq 2\|A'(\psi)\| \|\psi\|$$

$$\leq \epsilon \|A'(\psi)\|^2 + \frac{1}{\epsilon} \|\psi\|^2 \leq \epsilon (K \|\bar{\nabla} \psi\|^2 + K \|\psi\|^2) + \frac{1}{\epsilon} \|\psi\|^2$$

$$\Rightarrow 2\langle \Delta \psi, \psi \rangle \leq (2 + \epsilon K) \|\bar{\nabla} \psi\|^2 + (\epsilon K + \frac{1}{\epsilon}) \|\psi\|^2$$

$$\frac{1}{2} \cdot 2\langle \Delta \psi, \psi \rangle \leq \frac{2 + \epsilon K}{2} \|\bar{\nabla} \psi\|^2 + \frac{\epsilon K + \frac{1}{\epsilon}}{2} \|\psi\|^2$$

$$\Rightarrow \langle \Delta \psi, \psi \rangle \leq M (\|\bar{\nabla} \psi\|^2 + \|\psi\|^2) \\ \leq M (\|\nabla \psi\|^2 + \|\bar{\nabla} \psi\|^2 + \|\psi\|^2)$$

$$\Rightarrow \langle \Delta \psi, \psi \rangle + \langle \psi, \psi \rangle \leq M' (\|\nabla \psi\|^2 + \|\bar{\nabla} \psi\|^2 + \|\psi\|^2) \\ \text{" } \mathcal{Q}(\psi) \leq C \|\psi\|_1^2$$

This implies that the Dirichlet norm is equivalent to Sobolev 1-norm.

Remark 2. In the Kähler case, one may use the precise Weitzenböck formula, and the above integration by parts calculation to prove the Kodaira identity

$$\langle \Delta \psi, \psi \rangle = \|\bar{\nabla} \psi\|^2 + \langle R \psi, \psi \rangle$$

where, for $\psi \in A^{0,q}(M)$ and summing repeated indices,

$$\langle R \psi, \psi \rangle = q \int_M (R_{i\bar{j}} \psi_{\bar{c}_1 \dots \bar{c}_{q-1} \bar{i}} \bar{\psi}_{\bar{c}_1 \dots \bar{c}_{q-1} \bar{j}}) \Omega.$$